

Perturbation of the accretion flow onto supermassive black hole by a passing star

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- Supermassive black holes – surrounded by dust and gas (the constituents of the accretion flow), different stellar-size objects (stars, neutron stars or stellar-mass black holes) and (possibly) strong magnetic field
- Both the classical accretion discs or the radiative inefficient accretion flows (RIAF) can be perturbed by the stellar-size objects
- Interaction between different constituents has observable consequences \Rightarrow multiwavelength variability of galactic nuclei
- Can stars push out gas into jets (enrich jet by baryons) or form compact corona above the nuclei (lamp-post corona model)?
- GRMHD simulation of the flow \rightarrow implementation into HARM

- Open source software package for GRMHD computations HARMPI (Gammie et al, 2003; Tchekhovskoy et al., 2016)
- ideal MHD – no resistivity, magnetic field frozen into gas
- solver for continuity ($(\rho u^\mu)_{;\mu} = 0$) and energy-conservation equation ($T^\mu{}_{\nu;\mu} = 0$; $T_{gaz}^{\mu\nu} = (\rho + \rho\varepsilon + p)u^\mu u^\nu + pg^{\mu\nu}$)
- conservative scheme:

$$\partial_t \mathbf{U}(\mathbf{P}) = -\partial_i \mathbf{F}^i(\mathbf{P}) + \mathbf{S}(\mathbf{P}) \quad (1)$$

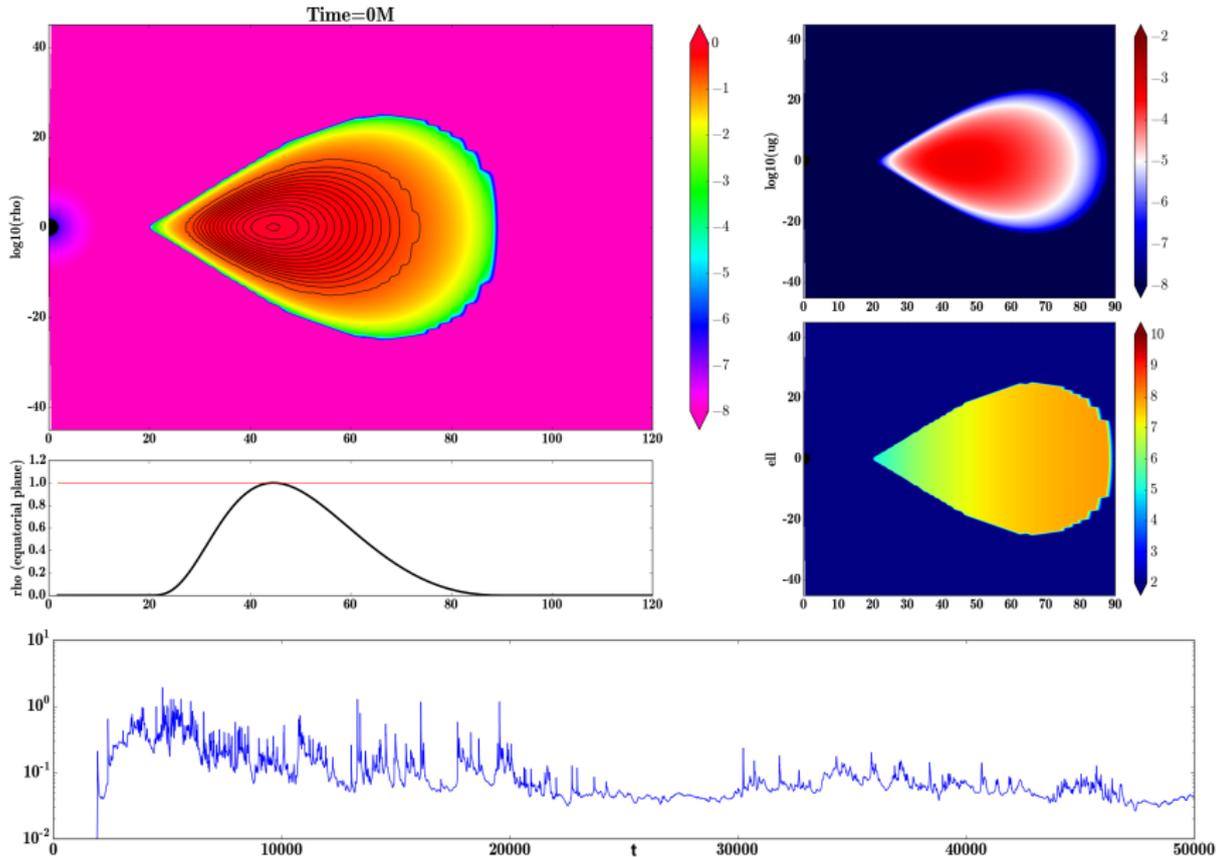
\mathbf{U} – conserved vars, \mathbf{P} – primitive vars, \mathbf{F}^i – fluxes, \mathbf{S} – sources

- numerical inversion of non-linear relation $\mathbf{U}(\mathbf{P})$
- fixed background (Kerr metric) – faster computation
- spherical coordinates – suitable for our geometry
- logarithmic grid in r – no need of grid refinement

Initial conditions – evolved thick torus (RIAF)

- Classical solutions – Fishbone & Moncrief, 1976; Abramowicz et al., 1978 – thick torus without magnetic field
- Komissarov, 2006, Montero et al., 2007
 - magnetized solution with purely azimuthal magnetic field
 - different EOS, constant angular momentum
- Witzany & Jefremov, 2018
 - extension to 2-parametric family of tori, non-constant profile of angular momentum (both growing and decreasing), closed form of the solution
 - different shapes of the torus (height versus radial extent)
- Our initial torus is large (20M – 90M) – reservoir of matter
- Instead of azimuthal magnetic field, the torus is interlaced with weak magnetic field – field lines inside the torus follow equipotentials of density \Rightarrow MRI in short time

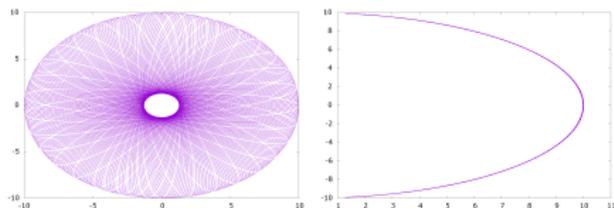
Initial conditions – evolved thick torus (RIAF)



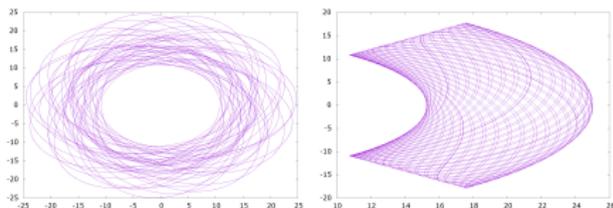
- Motion of a star in the accreting medium is complicated – approximations needed
- Dynamical effect of the star on the gas
 - Consecutive perturbation by moving star – velocity of gas in cells corresponding to the volume of the star is prescribed to equal star's velocity
- We do not evolve the stellar structure nor the feedback from the accretion flow on the star trajectory – star motion solved by geodesic equation on given Kerr background
- Evolved state of the torus ($t = 25000M$) as initial condition (ideally in a relaxed stationary accretion state)
- 2D runs: ϕ coordinate of the star is “forgotten”
- Changes in accretion rate + outflows

- 7x 2D runs: different star trajectory begins at $t = 25000M$ of the “basic” run
 $t_f = 50000M$ or $100000M$
star radius usually $r_{\text{star}} = 1M$ (for Sgr A* $1M \sim 9r_{\odot}$)
resolution 252×192
 - Run 1: Star moves close to $\sim 10M$, into “funnel” along axis
 - Run 2: Star moves between $\sim 12M - 25M$, not in funnel
 - Run 3: Star moves between $\sim 25M - 50M$, embedded in disc, close to equatorial plane
 - Run 4: Star moves close to $\sim 50M$, into funnel close to axis
 - Run 5: Star moves close to $\sim 50M$, not into funnel
 - Run 6: Star same as Run 5, radius of star $R = 10M$
 - Run 7: Star same as Run 1, radius of star $R = 0.1M$
- 1x 3D run – Run 8: Star same as Run 1, $t_f = 30000M$, resolution $252 \times 192 \times 96$

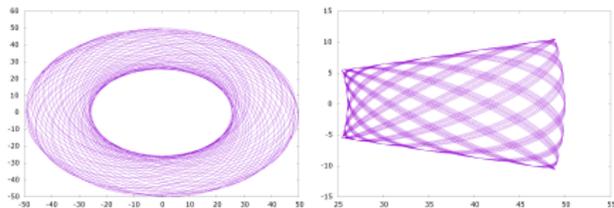
Run 1, 7



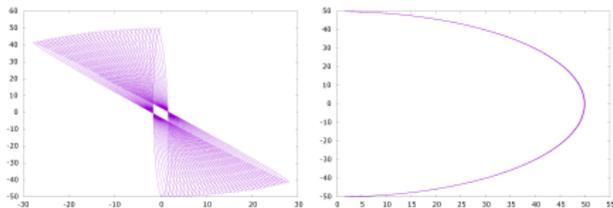
Run 2



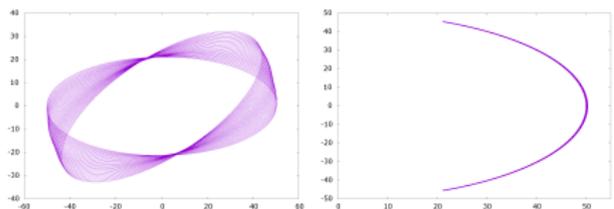
Run 3



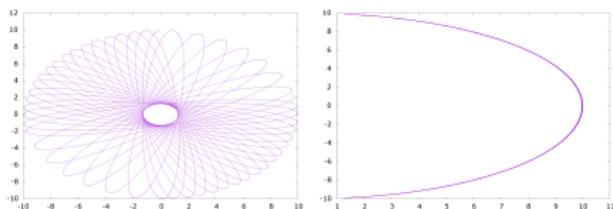
Run 4



Run 5, 6



Run 8



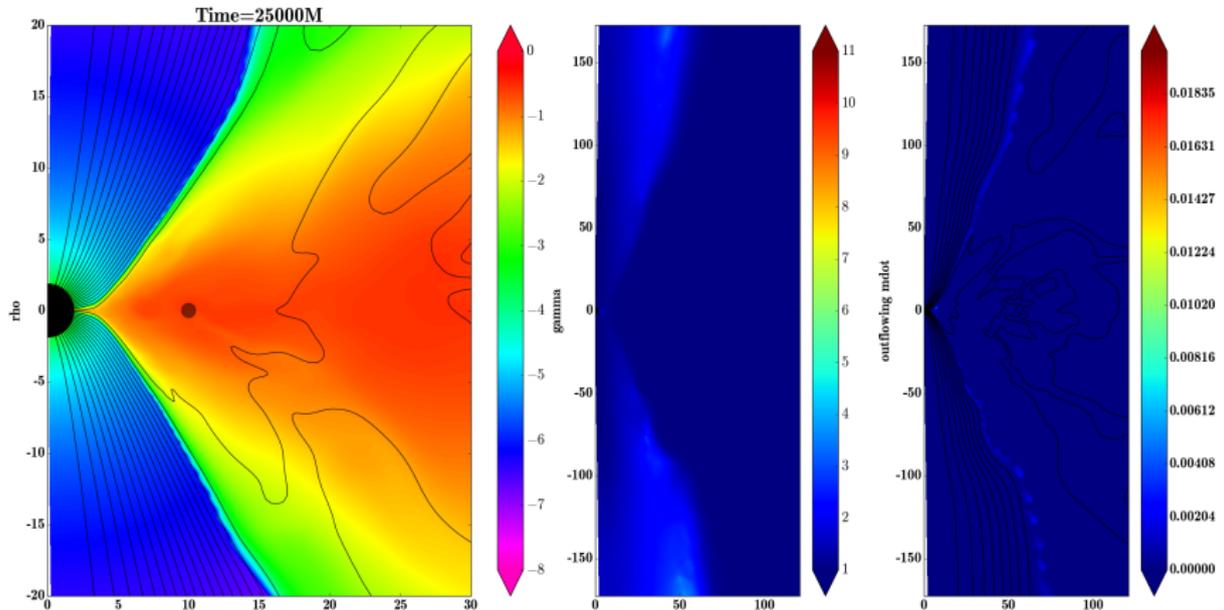
x-y plane

r-z plane

x-y plane

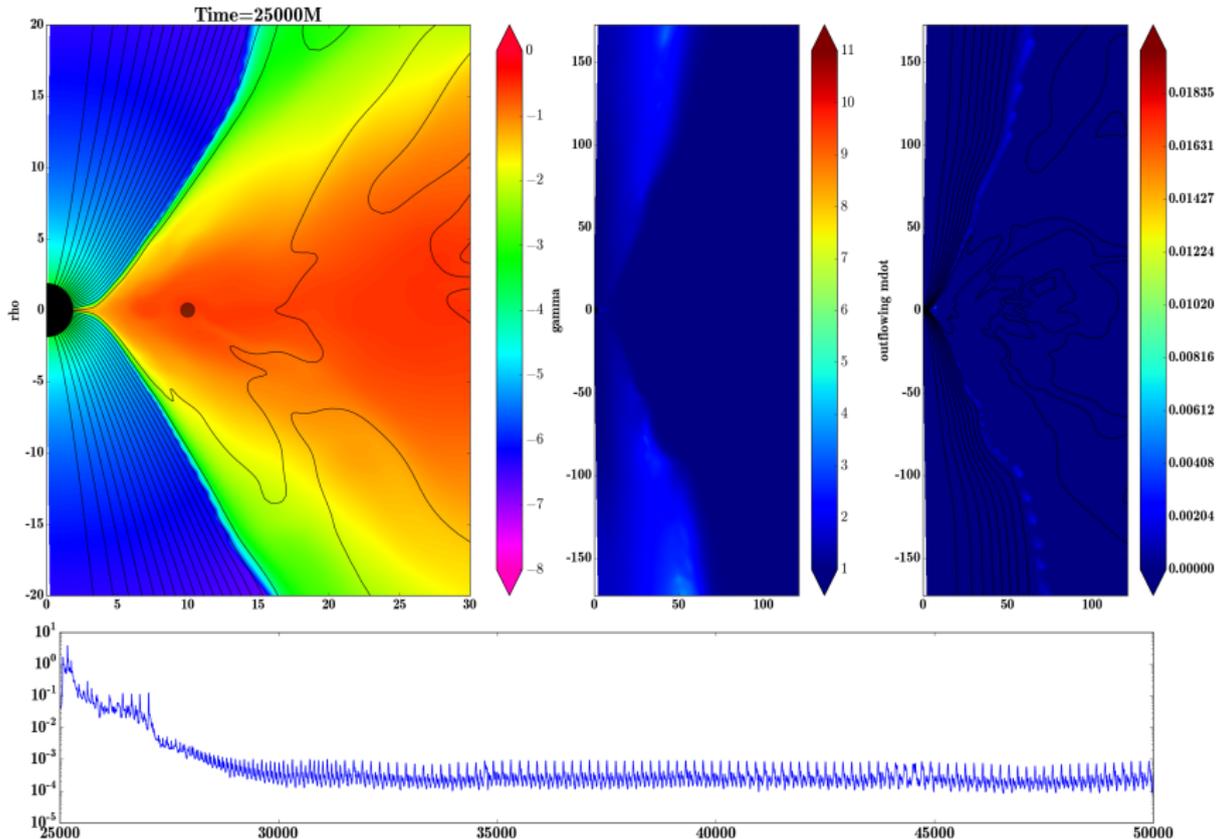
r-z plane

Run 1 – gas evolution

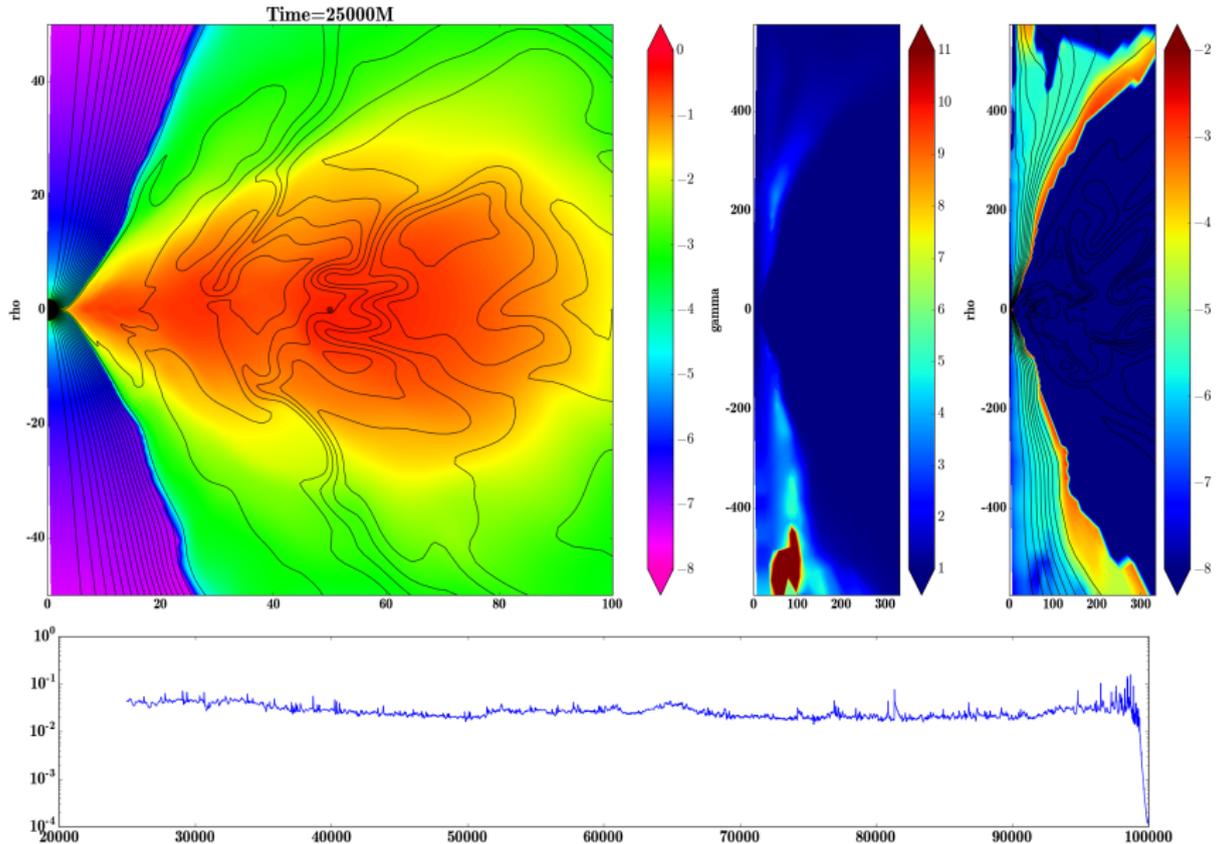


$$\dot{M}_{\text{out}}(r) = \rho(r) \frac{u^r(r)}{u^t(r)} \sqrt{-g} d\theta d\phi \quad \text{for } \gamma > 1.155$$

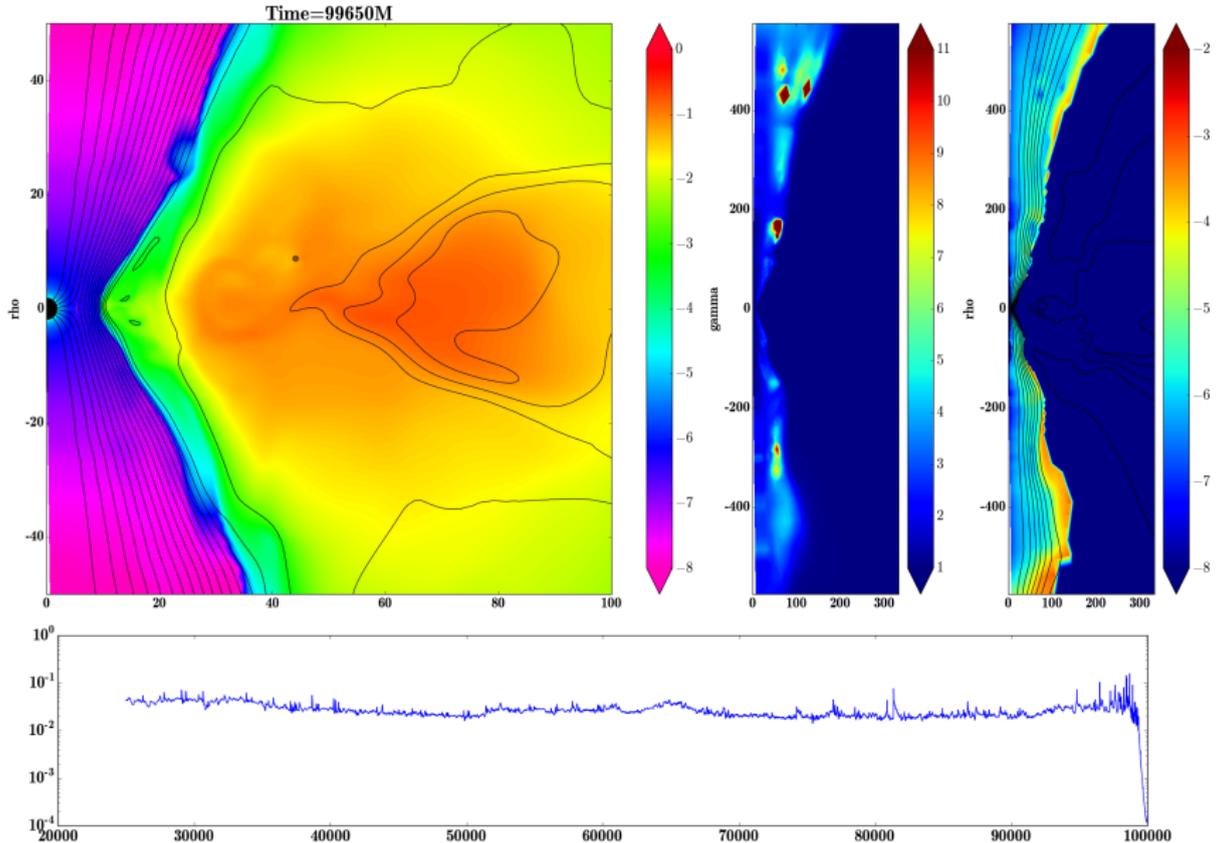
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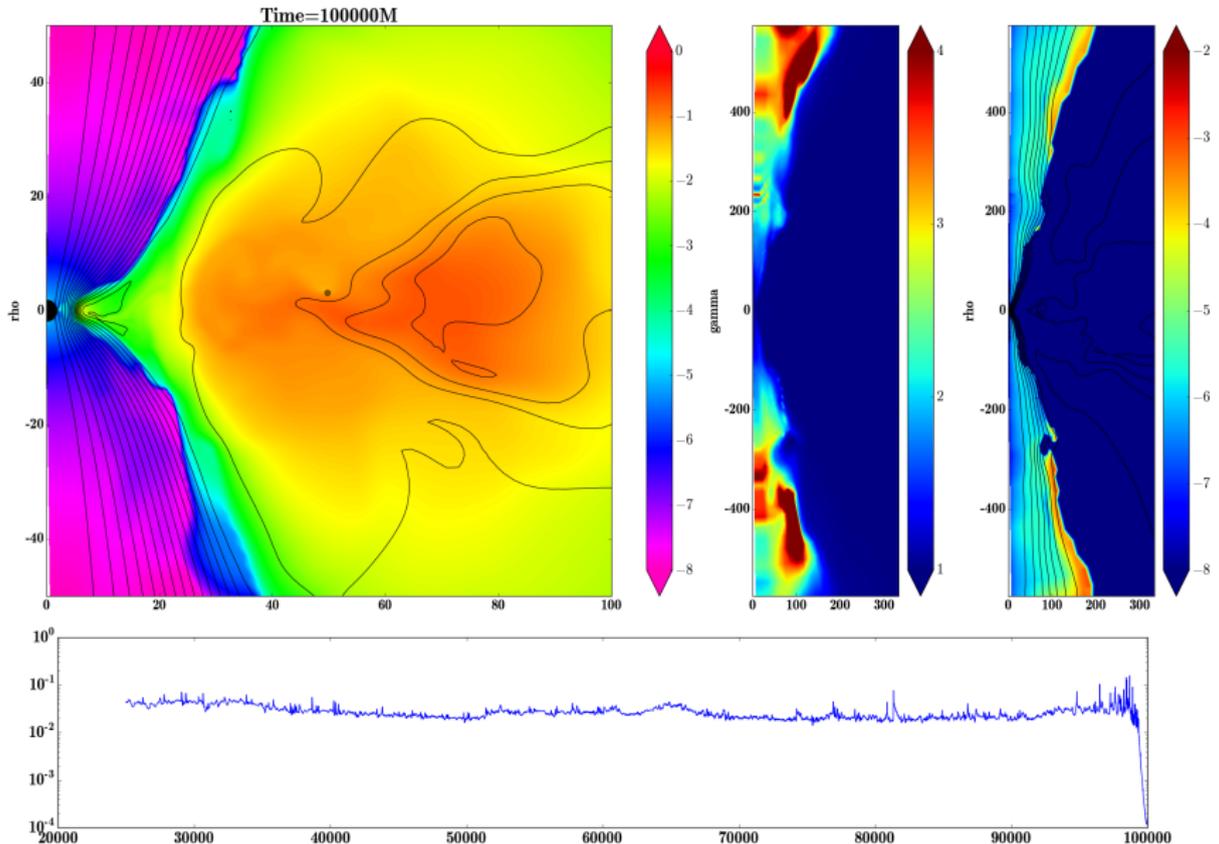
Run 3 – gas evolution



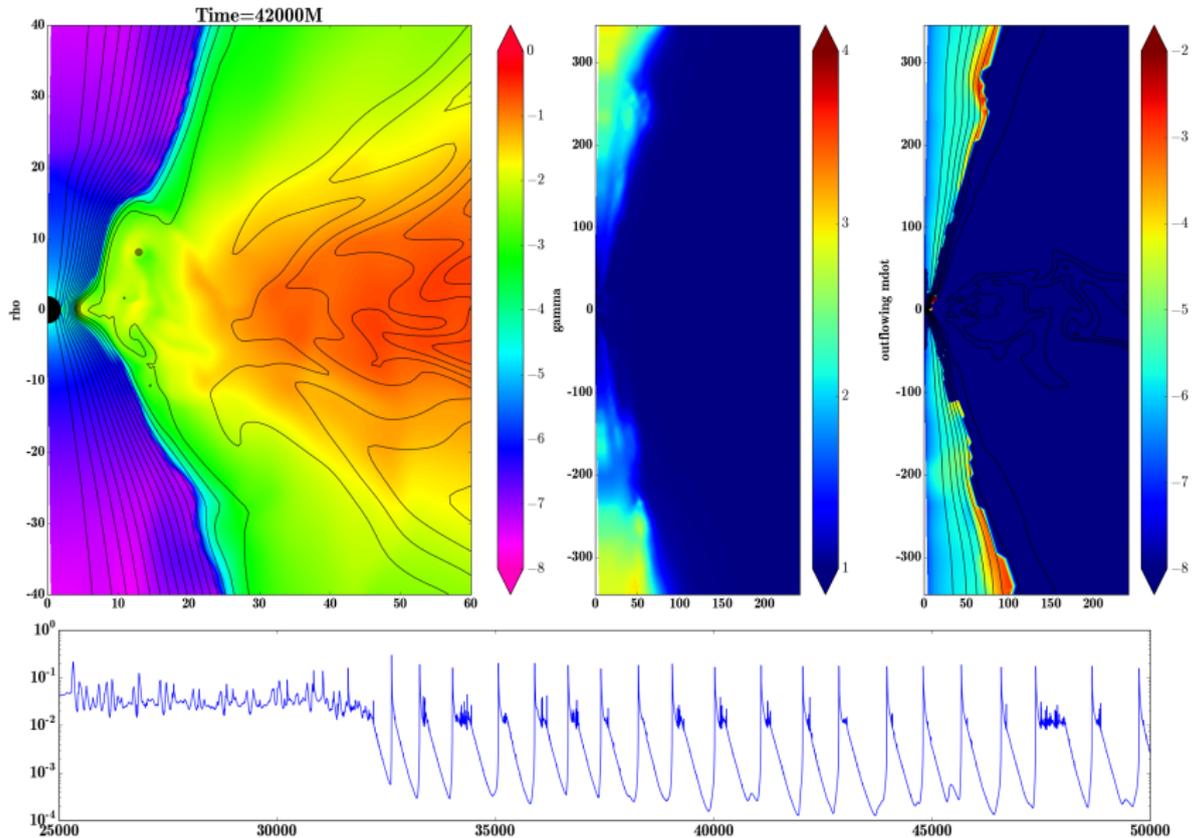
Run 3 – gas evolution



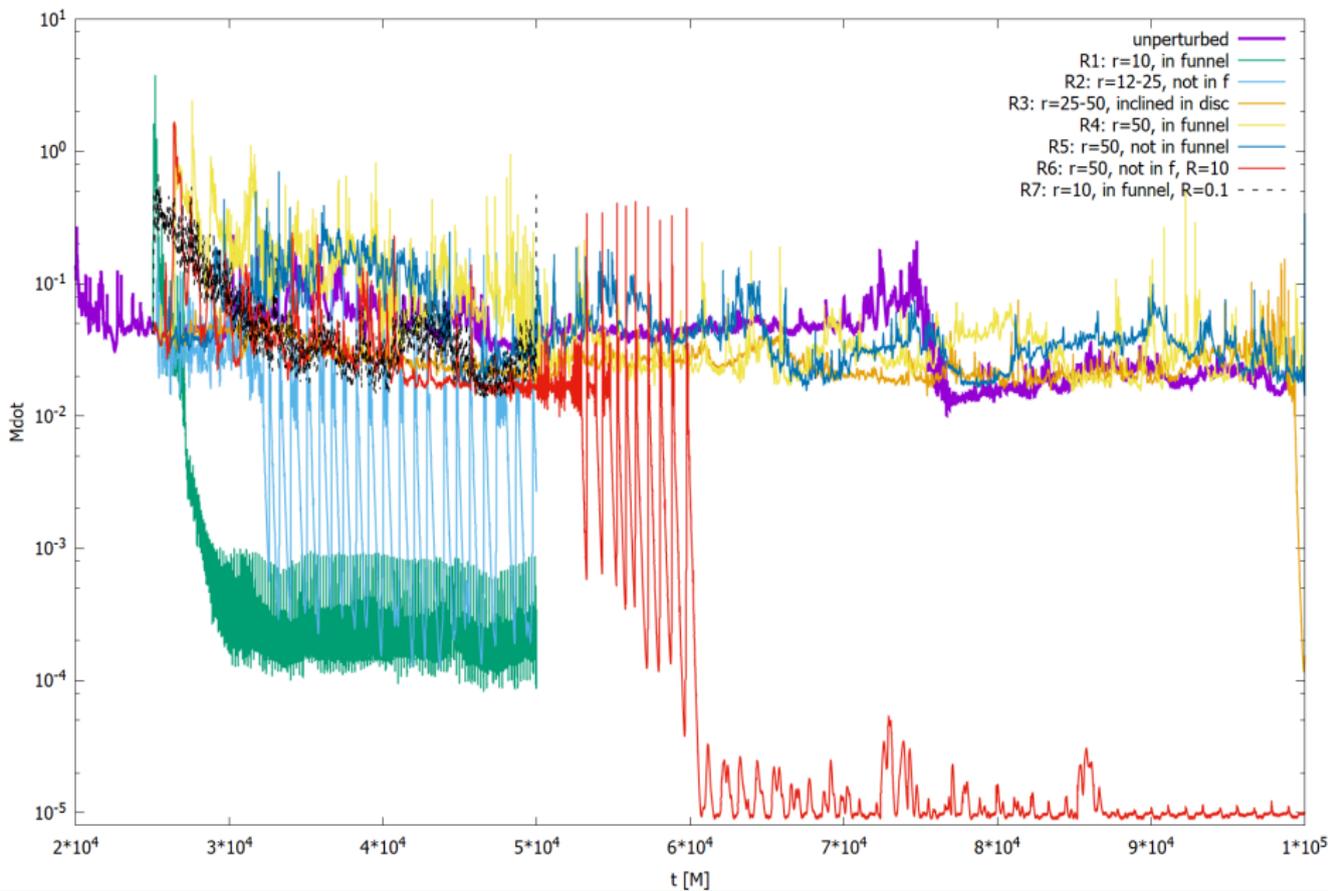
Run 3 – gas evolution



Run 2 – gas evolution

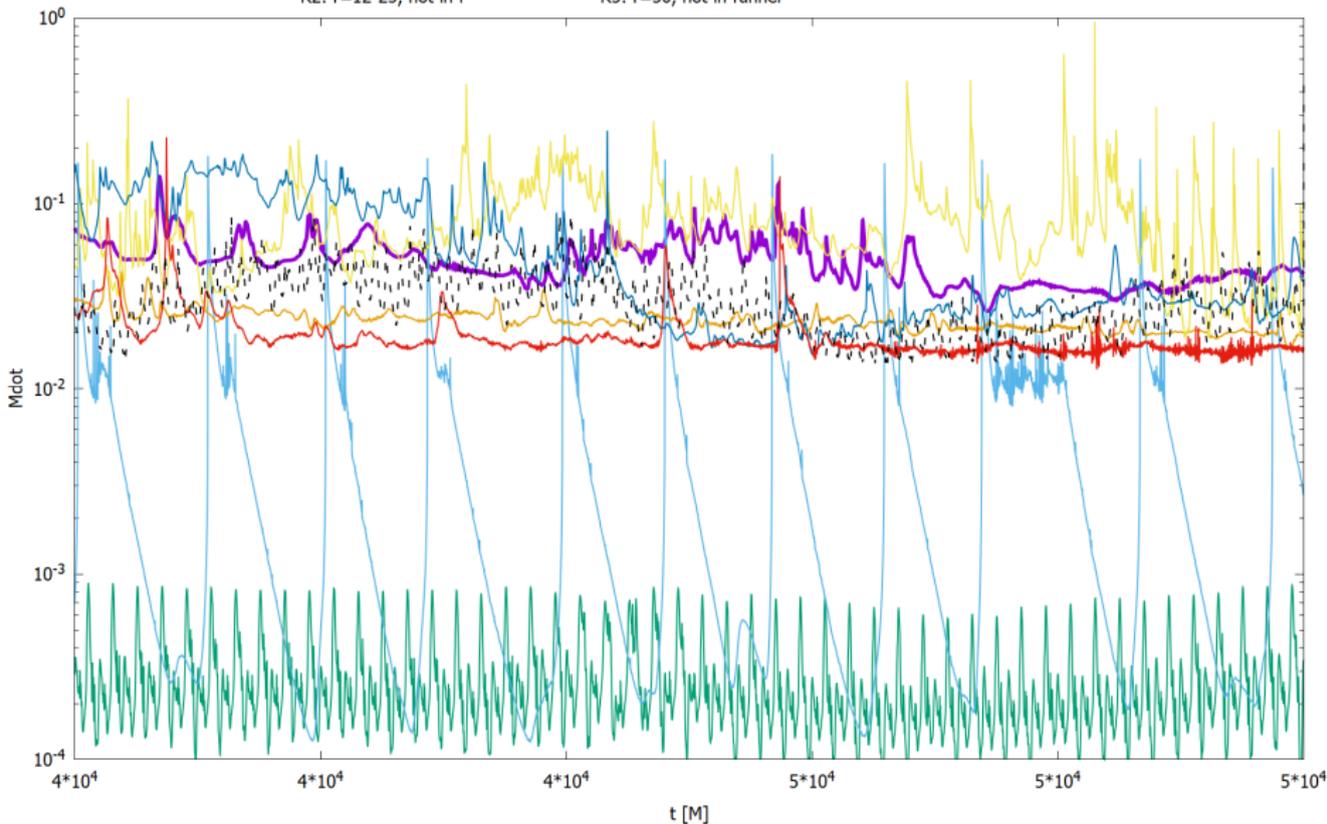


mdot of unperturbed run + 7 perturbed runs

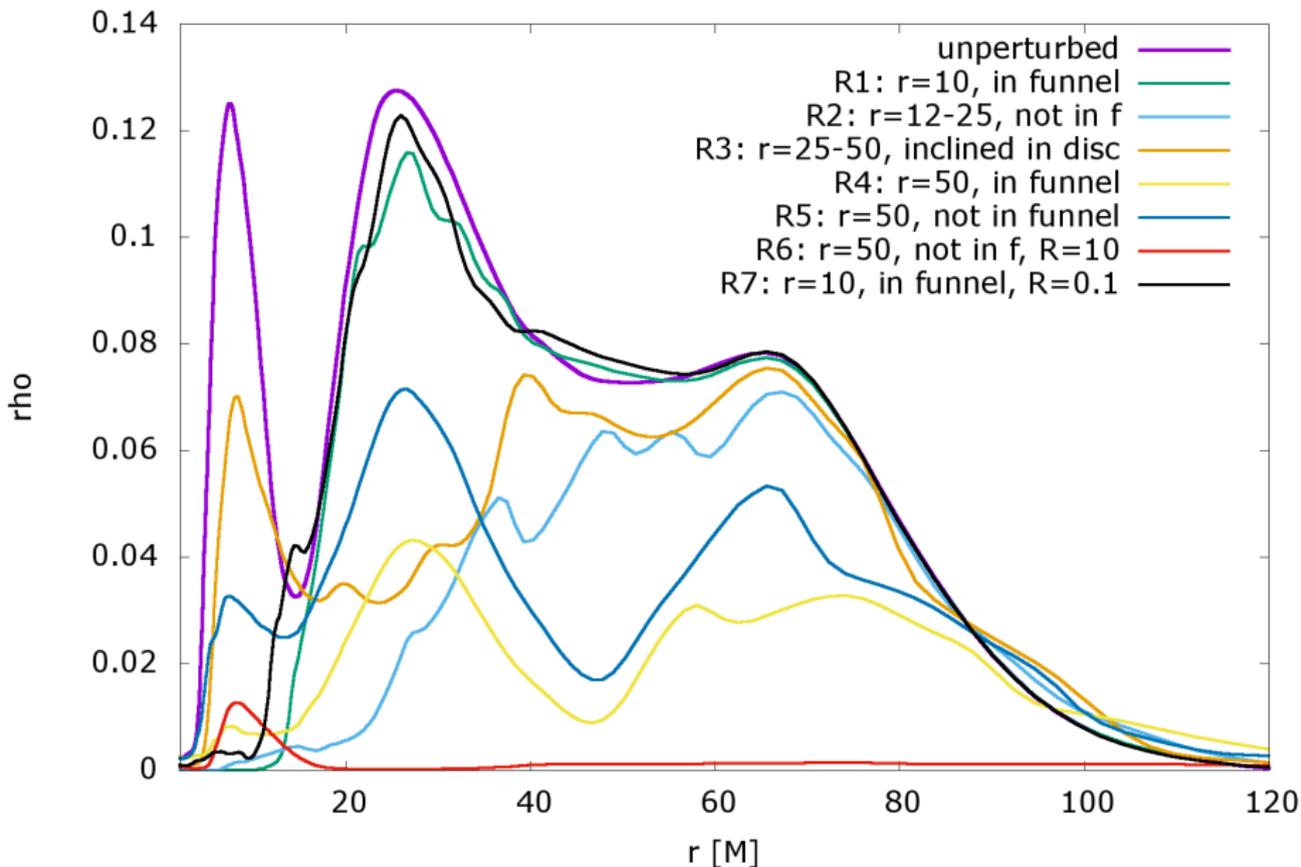


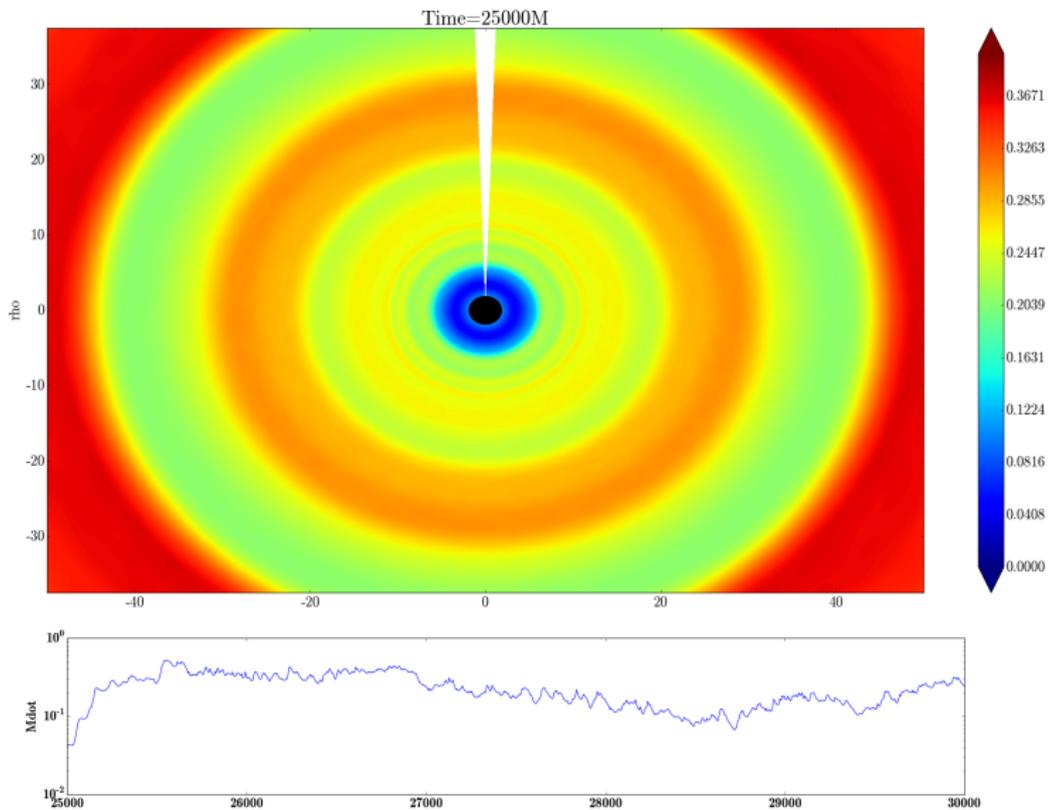
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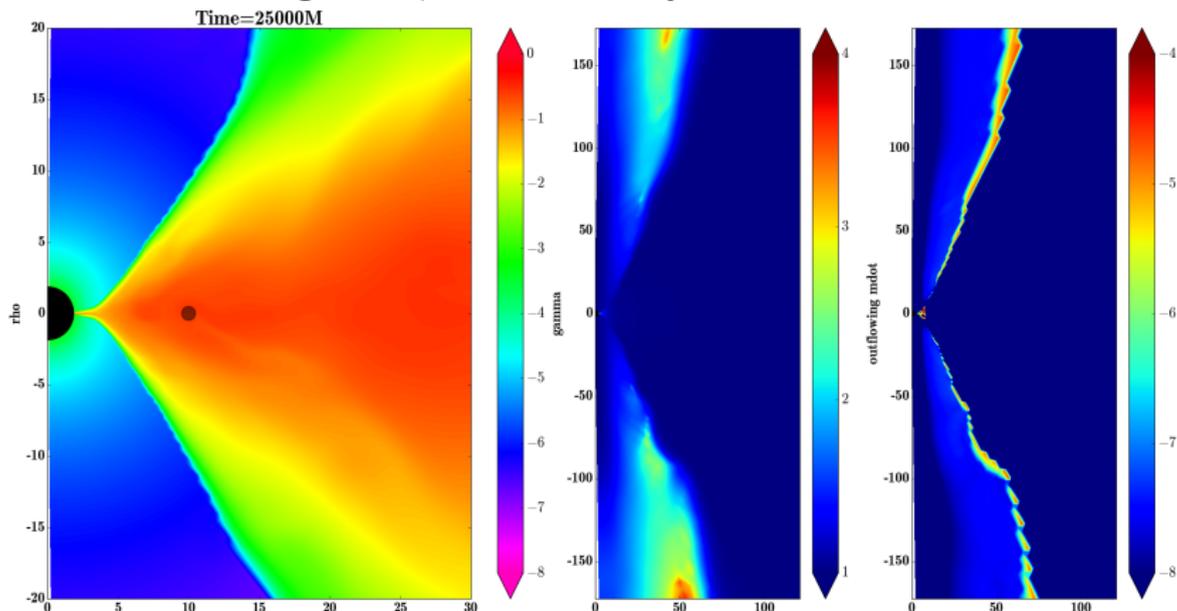
unperturbed — purple
R1: r=10, in funnel — green
R2: r=12-25, not in f — light blue
R3: r=25-50, in disc — orange
R4: r=50, in funnel — yellow
R5: r=50, not in funnel — blue
R6: r=50, not in f, R=10 — red
R7: r=10, in funnel, R=0.1 — black dashed



θ -averaged profile of ρ at $t = 50000M$



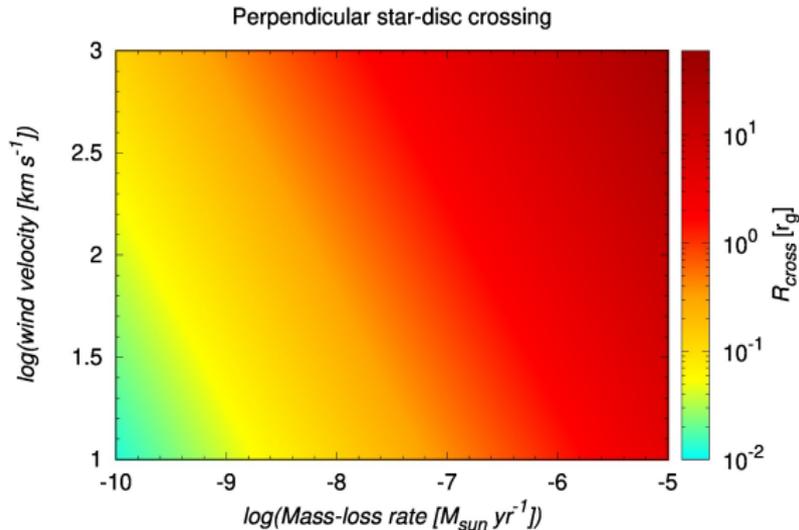


3D result averaged in ϕ -direction only for visualisation

3D run much shorter than 2D runs - stationary state not achieved
 similar properties – outflowing blobs along edge of funnel
 quantitative differences – smaller amount of outflowing mass

- Computation in geometrical units - scaled by black hole mass
- Interplay between BH mass, star's physical size and mass and the resolution of the grid
- Tidal radius r_t - distance from BH where the star rips
 - Sgr A*: $M_{\text{BH}} = 4.2 \cdot 10^6 M_{\odot} \rightarrow 1M = 6.3 \cdot 10^6 \text{ km}$
 - Sun: $1r_{\odot} = 7 \cdot 10^5 \text{ km} = 1/9M \rightarrow r_t = 18M$
 - massive star $M \sim 100M_{\odot}, r \sim 9r_{\odot} = 1M \rightarrow r_t = 35M$
 - white dwarf
 $M \sim 0.5M_{\odot}, r \sim 0.015r_{\odot} = 0.002M \rightarrow r_t = 0.33M$
- Grid resolution near $r \sim 10M \rightarrow \sim 0.2M$
Most of our runs – star radius $R = 1M$
test with $R = 0.1M$ – worked well, so far the lower limit

- Star with strong wind (S2: $M = 13.6M_{\odot}$, $r = 5.53r_{\odot} = 0.6M$)
- Stagnation radius R_{stag} = equilibrium between wind kinetic pressure and gas ram pressure
- Wilkin (1996): Cross-section of bow shock $r = 1.73R_{\text{stag}}$ (perpendicular crossing) – up to 10M at radii $r \sim 50M$



- Preliminary observations
 - Outgoing density waves in the accretion flow
 - Outgoing relativistic blobs in the funnel (γ up to several units)
 - the gaseous blobs go along the boundary between the funnel and the torus
 - Effects on the accretion rate :
 - decreased up to 2.5 orders of magnitude (4 orders in extreme case)
 - or even increased for some time period
 - quasiperiodic peaks and drops
 - stronger effect in 2D runs as anticipated
 - Influence on the total mass of the torus and the matter distribution
 - 3D simulation shows similar features with quantitatively lower effects

Thank you for your attention.